

LES LECTURE 23 (page 1)

Reminder papers discuss presentation schedule.

our exam is wed 4th 8-10 am

14 people to present \Rightarrow 2 class periods.

\hookrightarrow 120 marks

- Paper format

• 1st go over BC evaluations...

• PDF methods (rest of class will be this)

(~~and~~ we will mostly follow page ch. 12)

Review

Definitions for PDFs

PDF of vel. \nearrow in simple space V for velocity \vec{v}
at pos. \vec{x} and t

$$\int f(\vec{v}; \vec{x}, t) d\vec{v} = 1$$

• mean (or expected) value is defined by

$$\langle Q(\vec{v}; \vec{x}, t) \rangle = \int Q(\vec{v}) f(\vec{v}; \vec{x}, t) d\vec{v}$$

\Rightarrow mean velocity is simply.

$$\langle \vec{v}(\vec{x}, t) \rangle = \int \vec{v} f(\vec{v}; \vec{x}, t) d\vec{v}$$

and

$$\langle u_i u_j \rangle = \int (v_i - \langle u_i \rangle) (v_j - \langle u_j \rangle) f(v; \vec{x}, t) dv$$

LES LECTURE 23

page 2

in the PDF transport equation unknowns appear as conditional means.

$\phi(x, t) \equiv$ random field

$f_{U\phi} \equiv$ (one point one time) joint PDF of U and ϕ
sample space of ϕ
 $(V, \psi; \vec{x}, t)$

PDF of $\phi(x, t)$ conditional on $U(\vec{x}, t) = V$ is

$$f_{\phi|U}(\psi | \vec{V}, \vec{x}, t) = \frac{f_{U\phi}(\vec{V}, \psi; \vec{x}, t)}{f(\vec{V}; \vec{x}, t)} \quad (*)$$

and the conditional mean value is:

$$\langle \phi(\vec{x}, t) | U(\vec{x}, t) = V \rangle = \int \psi f_{\phi|U}(\psi | \vec{V}, \vec{x}, t) d\psi$$

may thus be written as

$$\langle \phi | \vec{V} \rangle$$

~~the~~ this conditional mean ~~can~~ be associated with the actual mean

$$\begin{aligned} \langle \phi(x, t) \rangle &= \iint \psi f_{U\phi}(\vec{V}, \psi; \vec{x}, t) d\psi d\vec{V} \\ &= \int \langle \phi(\vec{x}, t) | U(\vec{x}, t) = V \rangle f(\vec{V}; \vec{x}, t) d\vec{V} \\ &= \int \langle \phi | \vec{V} \rangle f d\vec{V} \quad (\text{in shorthand}) \end{aligned}$$

use (*)

PDF Transport equation (Pope appendix H)

• derivation relies on the idea of a fine grained PDF

$$f'(\vec{v}; \vec{x}, t) \equiv \delta(\vec{U}(\vec{x}, t) - \vec{v}) = \prod_{i=1}^3 \delta(U_i(\vec{x}, t) - v_i)$$

at ~~every~~ every point and time in the flow field f' is a delta function at $\vec{U}(\vec{x}, t) = \vec{v}$
 ↳ appendix C

• Two properties of fine grained PDF

$$\langle f'(\vec{v}; \vec{x}, t) \rangle = f(\vec{v}; \vec{x}, t) \quad \text{and}$$

$$\langle \phi(\vec{x}, t) f'(\vec{v}; \vec{x}, t) \rangle = \langle \phi(\vec{x}, t) | \vec{U}(\vec{x}, t) = \vec{v} \rangle f(\vec{v}; \vec{x}, t)$$

this comes from (do substitution relation)

$$\begin{aligned} \langle f'(\vec{v}; \vec{x}, t) \rangle &= \langle \delta(\vec{U}(\vec{x}, t) - \vec{v}) \rangle \\ &= \int \delta(\vec{v}' - \vec{v}) f(\vec{v}'; \vec{x}, t) d\vec{v}' \\ &= f(\vec{v}; \vec{x}, t) \end{aligned}$$

$$\int g(\vec{x}) \delta(\vec{x} - \vec{y}) = g(\vec{y})$$

shifting property

obtained in the same way (see Pope)

• derivatives of fine-grained PDFs.

using Pope's example: take $f'_a(\vec{v}; t) = \delta(u(t) - v)$
 sample-space variable.
 a constant
 scalar process $u(t)$

even $[f(x) = f(-x)]$

odd $[-f(x) = f(-x)]$

$$\frac{d}{dv} \delta(v-a) = \frac{d}{dv} \delta(a-v) \Rightarrow \delta^{(n)}(v-a) = -\delta^{(n)}(a-v)$$

(derivative is an odd function)

notation for deriv.

LES LECTURE 23 page 41

• differentiating our scalar from general PDF
Chain Rule ($\frac{dy}{dx} = \frac{dy}{du} \times \frac{du}{dx}$)

$$\begin{aligned} \frac{\partial}{\partial t} f'_u(v; t) &= g^{(1)}(u(t)-v) \frac{du(t)}{dt} = -g^{(1)}(v-u(t)) \frac{du(t)}{dt} \\ &= - \frac{\partial f'_u(v; t)}{\partial v} \frac{du(t)}{dt} = - \frac{\partial}{\partial v} \left(f'_u(v; t) \frac{du(t)}{dt} \right) \end{aligned}$$

• we also need that

$u(t)$ is independent of v

**

$$\begin{aligned} v_i(\vec{x}, t) \frac{\partial}{\partial x_i} f'(\vec{v}; \vec{x}, t) &= \frac{\partial}{\partial x_i} [v_i(\vec{x}, t) f'(\vec{v}; \vec{x}, t)] \\ &= \frac{\partial}{\partial x_i} [v_i f'(\vec{v}; \vec{x}, t)] = v_i \frac{\partial}{\partial x_i} f'(\vec{v}; \vec{x}, t) \end{aligned}$$

- uses incompress (to bring into dot)
- shifting property of δ functions
- v_i is independent (can move out)

Derivation of PDF Transport Eqn.

substantial (or total) derivative of f'

$$* \quad \frac{Df'}{Dt} = \frac{\partial f'}{\partial t} + v_i \frac{\partial f'}{\partial x_i} = - \frac{\partial}{\partial v_i} \left(f' \frac{Dv_i}{Dt} \right)$$

we get this using the equivalent of our scalar derivation of f' in time and space (for vector f')

* The main of this is

$$\frac{\partial f}{\partial t} + v_i \frac{\partial f}{\partial x_i} = - \frac{\partial}{\partial v_i} \left(f \left\langle \frac{Dv_i}{Dt} \mid \vec{v} \right\rangle \right) \quad \left(\text{for incompressible flow since } \nabla \cdot \vec{v} = 0 \text{ used} \right)$$

this is general and has no physics (we haven't said anything about the medium)

we insert physics by substituting in our N-S definition of $\frac{Dv_i}{Dt}$ (typical RHS of N-S)

LES LECTURE 23 page 5

with our N-S equations

$$\frac{\partial f}{\partial t} + v_i \frac{\partial f}{\partial x_i} = - \frac{\partial}{\partial v_i} \left(f \left\langle v \nabla^2 v_i - \frac{1}{\rho} \frac{\partial p}{\partial x_i} \right| \vec{v} \right\rangle \right)$$

if we decompose pressure as

$$\left\langle \frac{\partial p}{\partial x_i} \right| \vec{v} \right\rangle = \frac{\partial \langle p \rangle}{\partial x_i} + \left\langle \frac{\partial p'}{\partial x_i} \right| \vec{v} \right\rangle$$

we get

$$\frac{\partial f}{\partial t} + v_i \frac{\partial f}{\partial x_i} = \frac{1}{\rho} \frac{\partial \langle p \rangle}{\partial x_i} \frac{\partial f}{\partial v_i} - \frac{\partial}{\partial v_i} \left[f \left\langle v \nabla^2 v_i - \frac{1}{\rho} \frac{\partial p'}{\partial x_i} \right| \vec{v} \right\rangle \right]$$